One-Dimensional Diffusion Operators

Stanley Sawyer — Washington University — Vs. July 7, 2008

1. Basic Assumptions. Let s(x) be a continuous, strictly-increasing function s(x) on $I_0 = (0, 1)$ and let m(dx) be a Borel measure m(dx) on I_0 (that is, $m(K) < \infty$ for compact subsets of $K \subset I_0$). Then we can define a "Feller-type operator"

$$Lf(x) = \frac{d}{dm(x)} \frac{d}{ds(x)} f(x)$$
(1.1)

for appropriate functions f(x) on I_0 . Here s(x) is called the *scale function* of Lf(x) and m(dx) is the *speed measure*. (See e.g. Feller 1952, 1954, 1955, Ito and McKean 1965, Sawyer 1974.)

Assume further that

$$s(x) \in C^{2}(I_{0}) \cap C(I), \quad s'(x) > 0, \quad s(dx) = s'(x)dx$$

 $m(dx) = m(x)dx, \quad m(x) > 0, \quad m(x) \in C(I_{0})$
(1.2)

for 0 < x < 1 and I = [0, 1]. Since $s(x) \in C(I)$, s(0) and s(1) are finite. (By C(S) we mean the set of all continuous functions on S and by $C^k(S)$ for an open set S the set of all k-times-continuously differentiable functions on S.) Then

$$Lf(x) = \frac{1}{m(x)} \frac{d}{dx} \left(\frac{1}{s'(x)} f'(x) \right)$$
(1.3)

if $f(x) \in C^2(I_0)$. In the following we also assume that

$$\int_{0}^{1/2} (s(y) - s(0)) m(dy) < \infty, \qquad \int_{1/2}^{1} (s(1) - s(y)) m(dy) < \infty$$
(1.4)

Condition (1.4) is called the condition that the operator Lf(x) has "exit boundaries" at 0 and 1.

The motivating example that we have in mind is

$$s(x) = x$$
 and $m(dx) = \frac{2dx}{x(1-x)}$

and more generally

$$s(x) = \frac{1 - e^{-2\gamma x}}{2\gamma}$$
 and $m(dx) = \frac{2e^{2\gamma x}dx}{x(1 - x)}$ (1.5)

for $\gamma \neq 0$. We leave it as an exercise to show that (1.3) and (1.5) correspond to the differential operator

$$Lf(x) = (1/2)x(1-x)f''(x) + \gamma x(1-x)f'(x)$$

2. An Integral Equation for Lf(x). Define

$$g(x,y) = \frac{\left(s(x \wedge y) - s(0)\right)\left(s(1) - s(x \vee y)\right)}{s(1) - s(0)}$$
(2.1)
$$= \begin{cases} \frac{\left(s(1) - s(x)\right)\left(s(y) - s(0)\right)}{s(1) - s(0)} & 0 \le y \le x \le 1\\ \frac{\left(s(1) - s(y)\right)\left(s(x) - s(0)\right)}{s(1) - s(0)} & 0 \le x \le y \le 1 \end{cases}$$

where $x \wedge y = \min\{x, y\}$, and $x \vee y = \max\{x, y\}$. Then since s(x) is increasing

$$g(x,y) = g(y,x)$$

$$g(x,x) \le \min\{s(1) - s(x), \ s(x) - s(0)\} \le \frac{s(1) - s(0)}{4}$$

$$g(x,y) \le \min\{g(x,x), \ g(y,y)\}$$
(2.2)

Define the linear operator

$$Kf(x) = \int_{0}^{1} g(x, y) f(y) m(dy)$$

= $\frac{s(1) - s(x)}{s(1) - s(0)} \int_{0}^{x} (s(y) - s(0)) f(y) m(dy)$
+ $\frac{s(x) - s(0)}{s(1) - s(0)} \int_{x}^{1} (s(1) - s(y)) f(y) m(dy)$ (2.3)

for $f(x) \in L^2(dm) = \{ f(x) : \int_0^1 f(x)^2 m(dx) < \infty \}$. By Cauchy's inequality

$$\int_0^1 |g(y,y)h(y)| m(dy) \le \sqrt{\int_0^1 g(y,y)^2 m(dy)} \sqrt{\int_0^1 h(y)^2 m(dy)}$$

where $g(x,y) \leq g(y,y) \leq C$ and $\int_0^1 g(y,y)m(dy) < \infty$ by (1.4) and (2.2). Thus by dominated convergence each $Kf(x) \in C_0(I_0)$ where

$$C_0(I_0) = \{ h(x) \in C(I_0) : \lim_{x \to 0} h(x) = \lim_{x \to 1} h(x) = 0 \}$$
(2.4)

The following theorem shows that the integral operator Kf(x) in (2.3) is the inverse of the differential operator Lf(x) in (1.1) and (1.3) with zero boundary conditions (that is, $Lf \in C_0(I_0)$).

Theorem 2.1. If $f(x) \in L^2(dm)$, then h(x) = Kf(x) is the unique function $h(x) \in C^1(I_0) \cap C_0(I_0)$ such that h'(x) is absolutely continuous in I_0 and Lh(x) = -f(x) for m(dx)-almost every $x \in I_0$. If also $f \in C(I_0)$, then $h \in C^2(I_0)$ and $L_xh(x) = -f(x)$ for all $x \in I_0$.

Proof. Assume h(x) = Kf(x) in (2.3). Thus

$$h(x) = \frac{s(1) - s(x)}{s(1) - s(0)} \int_0^x (s(y) - s(0)) f(y) m(dy) + \frac{s(x) - s(0)}{s(1) - s(0)} \int_x^1 (s(1) - s(y)) f(y) m(dy)$$

As above, $h \in C(I)$ with h(0) = h(1) = 0. Also, h(x) is absolutely continuous in x with

$$h'(x) = s'(x) \left(\int_x^1 \frac{s(1) - s(y)}{s(1) - s(0)} f(y) m(dy) - \int_0^x \frac{s(y) - s(0)}{s(1) - s(0)} f(y) m(dy) \right)$$
(2.5)

for almost every x. Since the right-hand side above is continuous, h'(x) extends to a continuous function on $C(I_0)$ such that both h'(x) and h'(x)/s'(x) are themselves absolutely continuous. Moreover, (d/dx)(h'(x)/s'(x)) = -f(x)m(x) for almost every x. Thus by (1.3) Lh(x) = -f(x) for m(dx)-almost every $x \in I_0$. If $f \in C(I_0)$, then $h \in C^2(I_0)$ and $L_x h = -f$ by the same proof.

Conversely, suppose that $h_1(x) \in C^1(I_0) \cap C_0(I_0)$, $h'_1(x)$ is absolutely continuous, and $Lh_1(x) = -f(x)$ for almost every x. I now claim that $h_1 = h = Kf(x)$. First, let $g(x) = h(x) - h_1(x)$. Then, following the same convention as in (2.5), $g(x) \in C^1(I_0) \cap C_0(I_0)$, g'(x) is absolutely continuous, and $L_xg(x) = Lh(x) - Lh_1(x) = 0$.

In generation, the relations (1.3) and $g \in C^1(I_0)$ with g'(x) absolutely continuous implies

$$g'(x) = g'(x_0) + s'(x) \int_{x_0}^x L_x g(y) m(y) dy$$
(2.6)

for $0 < x, x_0 < 1$. Thus g'(x) = 0 if $L_x g(x) = 0$ for m(dx)-a.e. x and g(x) is constant. Since g(0) = g(1) = 0, we must have $g = h - h_1 = 0$ and $h = h_1$, which completes the proof of Theorem 2.1.

We have actually proved a little more:

Lemma 2.1. Assume $g(x) \in C^1(I_0) \cap C(I)$, g(0) = g(1) = 0, g'(x) is absolutely continuous, and

$$Lg(x) = \frac{1}{m(x)} \frac{d}{dx} \left(\frac{g'(x)}{s'(x)} \right) = -\lambda g(x)$$
(2.7)

for m(dx)-a.e. x and some constant λ . Then either g(x) = 0 for 0 < x < 1 or else $\lambda > 0$.

Proof. By (1.3), (2.6) holds with $L_x g(y) = -\lambda g(y)$ for arbitrary $0 < x, x_0 < 1$. If $\lambda = 0$, then g'(x) = 0 for a.e. x and g(x) = g(0) = g(1) = 0. Now assume $\lambda \neq 0$.

Assume $\max_{0 \le x \le 1} g(x) = g(x_0) > 0$ and $\lambda < 0$. Then (2.6) implies that g'(x) > 0 for $x_0 < x < x_0 + \epsilon$ for some $\epsilon > 0$, which is inconsistent with the assumption that x_0 is an interior maximum of g(x). Thus $\lambda > 0$.

Lemma 2.2. Assume $f(x) \in L^2(I, dm)$ and

$$Kf(x) = \int_0^1 g(x, y)f(y) m(dy) = \lambda f(x) \text{ a.e. } dm$$
 (2.8)

for some constant λ . Then either f(x) = 0 a.e. dm. or else $\lambda > 0$.

Proof. By Theorem 2.1 and Lemma 2.1.

3. Kf(x) is a Hilbert-Schmidt Operator. That is, I claim Theorem 3.1. For g(x, y) in (2.1)

$$\int_{0}^{1} \int_{0}^{1} g(x,y)^{2} m(dx)m(dy) < \infty$$
(3.1)

Proof. By (2.2) and (1.4)

$$k(x) = \int_0^1 g(x, y)^2 \, m(dy) \, \le \, g(x, x) \, \int_0^1 g(y, y) \, m(dy) \, < \, \infty \tag{3.2}$$

Thus

$$\int_{0}^{1} \int_{0}^{1} g(x, y)^{2} m(dx) m(dy) = \int_{0}^{1} k(x) m(dx)$$
$$\leq \left(\int_{0}^{1} g(y, y) m(dy) \right)^{2} < \infty$$

We can use a similar argument to obtain

Theorem 3.2. Suppose $f_n, f \in L^2(I, dm)$ and

$$\lim_{n \to \infty} \int_0^1 \left(f_n(y) - f(y) \right)^2 m(dy) = 0$$

and Kf(x) is given by (2.3). Then

 $\lim_{n \to \infty} K f_n(x) = K(x) \quad \text{uniformly in } x$

Proof. Then

$$|Kf_n(x) - Kf(x)|^2 = \left| \int_0^1 g(x, y) (f_n(y) - f(y)) m(dy) \right|^2$$

$$\leq \int_0^1 g(x, y)^2 m(dy) \int_0^1 (f_n(y) - f(y))^2 m(dy)$$

$$= k(x) \int_0^1 (f_n(y) - f(y))^2 m(dy) \to 0$$

uniformly in x since k(x) is bounded by (3.2)–(3.2).

4. Eigenfunctions and Eigenvalues of Kf(x). Since Kf(x) is a (symmetric) Hilbert-Schmidt operator, or equivalently since g(x, y) is a Hilbert-Schmidt kernel in $L^2(I, dm)$ by Theorem 3.1, there exist eigenfunctions $\phi_n(x)$ of K in $L^2(I, dm)$

$$K\phi_n(x) = \int_0^1 g(x, y)\phi_n(y) \, m(dy) = \mu_n \phi_n(x) \tag{4.1}$$

for nonzero real eigenvalues $\mu_n \neq 0$ satisfying the following conditions (see e.g. Riesz and Nagy, 1955, p242). The functions $\phi_n(x)$ are orthonormal, that is

$$(\phi_n, \phi_m) = \int_0^1 \phi_n(x)\phi_m(x) \, m(dx) = \delta_{mn}$$
 (4.2)

where

$$(f,g) = \int_0^1 f(x)g(x) m(dx)$$
(4.3)

for $f, g \in L^2(I, dm)$, and any $f \in L^2(I, dm)$ can be written

$$f = g + \sum_{n=1}^{\infty} c_n \phi_n \tag{4.4}$$

where Kg = 0, $(g, \phi_n) = 0$ for all n, and the series converges in $L^2(I, dm)$ (Riesz and Nagy ibid.).

By Lemma 2.2 with $\lambda = 0$, Kg = 0 implies g = 0, so that g = 0in (4.4). Thus the functions $\phi_n(x)$ are complete orthonormal in $L^2(I, dm)$ and $c_n = (f, \phi_n)$ in (4.4). That is, the expansion (4.4) with g = 0 holds for any $f \in L^2(I, dm)$ with $c_n = (f, \phi_n)$. Similarly, the eigenvalues $\mu_n > 0$ by Lemma 2.2.

By Theorem 2.1, the relations (4.1)–(4.2) implies $\phi_n(x) \in C^2(I_0) \cap C_0(I_0)$, $\phi_n(0) = \phi_n(1) = 0$, and $\phi_n(x) = -\mu_n L \phi_n(x)$. Thus (4.1) is equivalent to

$$L_x \phi_n(x) = -\lambda_n \phi_n(x), \quad \lambda_n = 1/\mu_n, \quad \phi_n(0) = \phi_n(1) = 0$$
 (4.5)

We can use the Hilbert-Schmidt property (3.1) and the properties of $\phi_n(x)$ to find out more information about the eigenvalues μ_n .

Theorem 4.1. The function g(x, y) in (2.1) satisfies

$$g(x,y) = \sum_{n=1}^{\infty} \mu_n \phi_n(x) \phi_n(y)$$
(4.6)

with convergence in $L^2(I^2, dm^2)$. Moreover

$$\int_0^1 \int_0^1 g(x,y)^2 \, m(dx) m(dy) = \sum_{n=1}^\infty \mu_n^2 < \infty$$
(4.7)

Proof of Theorem 4.1. Since the functions $\phi_n(x)$ in (4.1) or (4.5) are complete orthonormal in $L^2(I_0, dm)$, the functions $\{\phi_m(x)\phi_n(y)\}$ are complete orthonormal in $L^2(I^2, dm^2)$ where $I^2 = I_0 \times I_0$ and $dm^2(dxdy) = m(dx)m(dy)$. This means that if k(x, y) is any measureable function on the unit square with $\int_0^1 \int_0^1 k(x, y)^2 m(dx)m(dy) < \infty$, then

$$k(x,y) = \sum_{n=1}^{\infty} \sum_{m=1}^{\infty} c_{mn} \phi_m(x) \phi_n(y)$$
 (4.8)

with convergence in $L^2(I^2, dm^2)$ and

$$c_{mn} = \int_0^1 \int_0^1 g(w, z) \phi_m(z) \phi_n(w) m(dz) m(dw)$$
(4.9)

Moreover

$$\sum_{n=1}^{\infty} \sum_{m=1}^{\infty} c_{mn}^2 = \int_0^1 \int_0^1 k(x, y)^2 \, m(dx) m(dy) < \infty$$
 (4.10)

Since g(x, y) = k(x, y) is in $L^2(I^2, dm^2)$ by (3.1), g(x, y) satisfies (4.8) in the same sense. Since $\int g(x, y)\phi_n(y)m(dy) = \mu_n\phi_n(x)$ and $(\phi_n, \phi_m) = \delta_{mn}$, the Fourier coefficients in (4.9) are $c_{mn} = \mu_n\delta_{mn}$. This implies (4.6) and (4.7).

Theorem 4.2. Under the assumptions (1.1)-(1.4)

$$|\phi_n(x)| \leq \lambda_n C_1 \sqrt{g(x,x)} \tag{4.11}$$

for an absolute constant C_1 . If in addition $\int_0^1 \sqrt{g(x,x)} m(dx) < \infty$,

$$|\phi_n(x)| \leq C_2 \lambda_n^2 g(x, x) \tag{4.12}$$

Proof. By (4.1) and Cauchy's inequality

$$\begin{aligned} |\phi_n(x)| &\leq \lambda_n \sqrt{\int_0^1 g(x, y)^2 m(dy)} \sqrt{\int_0^1 \phi_n(y)^2 m(dy)} \\ &\leq \lambda_n \sqrt{g(x, x)} \sqrt{\int_0^1 g(y, y) m(dy)} \end{aligned}$$
(4.13)

since $g(x, y) \leq \min\{g(x, x), g(y, y)\}$ by (2.2). The right-hand side of (4.13) is finite by (1.4). Similarly by (4.1) and (4.11)

$$\begin{aligned} |\phi_n(x)| &\leq \lambda_n \int_0^1 g(x,y) |\phi_n(y)| m(dy) \\ &\leq C_1 \lambda_n^2 g(x,x) \int_0^1 \sqrt{g(y,y)} m(dy) \end{aligned}$$

by (2.2).

We also have

Theorem 4.3. The series

$$g(x,y) = \sum_{n=1}^{\infty} \mu_n \phi_n(x) \phi_n(y)$$
 (4.14)

converges uniformly absolutely for $0 \le x, y \le 1$. In particular by (4.14)

$$\int_{0}^{1} g(x,x)m(dx) = \sum_{n=1}^{\infty} \mu_n \int_{0}^{1} \phi_n(x)^2 m(dx) = \sum_{n=1}^{\infty} \mu_n < \infty$$
 (4.15)

Proof. This is a special case of Mercer's Theorem (Riesz and Nagy, 1955, p245). We give the proof for completeness. First, consider the remainder

$$g_n(x,y) = g(x,y) - \sum_{k=1}^n \mu_k \phi_k(x) \phi_k(y) = \sum_{k=n+1}^\infty \mu_k \phi_k(x) \phi_k(y) \quad (4.16)$$

with convergence in $L^2(I^2, dm^2)$. Since the $\phi_n(x) \in C_0(I)$, the function $g_n(x,y)$ is continuous in I^2 and satisfies

$$\int_0^1 \int_0^1 f(x)f(y)g_n(x,y)m(dx)m(dy) = \sum_{k=n+1}^\infty \mu_k(f,\phi_k)^2 \ge 0 \quad (4.17)$$

by (4.16). I claim that this implies $g_n(x,x) \ge 0$ for all $n \ge 0$ and x. For, if $g_n(x_0, x_0) < 0$ for some $x_0 \in I_0$, then $g_n(x, y) < 0$ for $x, y \in (x_0 - \epsilon, x_0 + \epsilon)$ for some $\epsilon > 0$. Let $f(x) \in L^2(I, dm)$ be any positive continuous function with support in the interval $(x_0 - \epsilon, x_0 + \epsilon)$. Then

$$\int_0^1 \int_0^1 f(x)f(y)g_n(x,y)m(dx)m(dy)$$

=
$$\int_{x_0-\epsilon}^{x_0+\epsilon} \int_{x_0-\epsilon}^{x_0+\epsilon} f(x)f(y)g_n(x,y)m(dx)m(dy) < 0$$

which contradicts (4.17). It follows that $g_n(x, x) \ge 0$ and

$$0 \leq \sum_{k=1}^{n} \mu_k \phi_k(x)^2 \leq g(x, x)$$
(4.18)

for $0 \le x \le 1$. Next, by Cauchy's inequality

$$\sum_{k=m}^{n} |\mu_k \phi_k(x) \phi_k(y)| \leq \sqrt{\sum_{k=m}^{n} \mu_k \phi_k(x)^2} \sqrt{\sum_{k=m}^{n} \mu_k \phi_k(y)^2}$$
$$\leq \left(\sqrt{\sum_{k=m}^{n} \mu_k \phi_k(x)^2}\right) \sqrt{g(y,y)}$$
(4.19)

by (4.18). This implies that, for each fixed x, the series

$$B(x,y) = \sum_{n=1}^{\infty} \mu_n \phi_n(x) \phi_n(y)$$

converges uniformly in y. Recall that B(x, y) = g(x, y) a.e. dm^2 in the square by Theorem 4.1, but we will need pointwise identity.

For any $f(y) \in C_c(I_0)$ (that is, a continuous function of compact support $K \subset I_0$), the series

$$Bf(x) = \int_0^1 B(x, y) f(y) m(dy) = \sum_{n=1}^\infty \mu_n \phi_n(x) (f, \phi_n)$$

converges numerically for each x. By Theorem 3.2

$$Kf(x) = \int_0^1 g(x, y) f(y) m(dy) = \sum_{n=1}^\infty \mu_n \phi_n(x) (f, \phi_n)$$

with uniform convergence in x. It follows that

$$\int_0^1 (B(x,y) - g(x,y)) f(y) m(dy) = 0$$

for every x and $f \in C_c(I_0)$. Since we can take $f(y) = (B(x, y) - g(x, y))\psi(y)$ where $\psi(y) \ge 0$ and $\psi \in C_c(I_0)$, we conclude B(x, y) = g(x, y) for all (x, y). Thus

$$g(x,x) = B(x,x) = \sum_{n=1}^{\infty} \mu_n \phi_n(x)^2$$
 (4.20)

for all x with 0 < x < 1, with (4.20) being trivial for x = 0 and x = 1. Since both g(x, x) and $\phi_n(x)$ are continuous on the closed interval [0, 1], and since the convergence in (4.20) is monotonic since $\mu_n > 0$ and $\phi_n(x)^2 \ge 0$, it follows from Dini's theorem that the convergence in (4.20) is uniform in x. It then follows from (4.19) that the series (4.14) converges uniformly in I^2 , which completes the proof of the theorem.

(A quick proof of Dini's Theorem: For each $x \in [0, 1]$, the series converges within $2\epsilon > 0$ for $n \ge n_x$, and hence within $\epsilon > 0$ for $n \ge n_x$ and $|y - x| < \delta_x$ and $y \in I$ for some $\delta_x > 0$. Note $[0, 1] \subset \bigcup_x (x - \delta_x, x + \delta_x)$ and use the Heine-Borel theorem.)

5. Diffusion Semi-Groups and Non-negativity. Define

$$p(t,x,y) = \sum_{n=1}^{\infty} e^{-\lambda_n t} \phi_n(x) \phi_n(y)$$
(5.1)

for $0 < t < \infty$ for the eigenvectors and eigenvalues $\phi_n(x)$ and λ_n defined in (4.5) and (4.1). Since $|\phi_n(x)| \leq C_1 \lambda_n$ by (4.11) where $\sum_{n=1}^{\infty} (1/\lambda_n) < \infty$, the

series converges uniformly for $0 \le x, y \le 1$ and $a \le t < \infty$ for any a > 0. By Mercer's Theorem (Theorem 4.3), $\sum (1/\lambda_n)\phi_n(x)\phi_n(y)$ converges absolutely uniformly for $0 \le x, y \le 1$. Thus

$$\int_0^\infty p(t, x, y) dt = g(x, y) \tag{5.2}$$

with uniform convergence of the integral for $0 \le x, y \le 1$.

Since the series (5.1) converges in $L^2(I, dm)$ for each x,

$$p(t+s,x,z) = \int_0^1 p(t,x,y)p(s,y,z)m(dy)$$
(5.3)

for t, s > 0. Similarly, define

$$Q_t f(x) = \int_0^1 p(t, x, y) f(y) m(dy), \qquad f \in L^2(I, dm)$$
(5.4)

It follows from (5.3) that

$$Q_{t+s}f(x) = Q_t(Q_s f)(x), \qquad s, t > 0$$
 (5.5)

The main result of this section is

Theorem 5.1. For p(t, x, y) in (5.1) and t > 0, $0 \le x \le 1$,

- (i) $p(t, x, y) \ge 0$ $(0 \le y \le 1)$
- (ii) $\int_0^1 p(t,x,y) m(dy) \leq 1$
- (iii) For any $f \in C_0(I_0)$,

$$\lim_{t \to 0} \int_0^1 p(t, x, y) f(y) m(dy) = f(x)$$
(5.6)

uniformly for $0 \le x \le 1$.

Proof. Choose $h(x) \in C^2(I)$ such that, for some a > 0, h(x) = 0 for $0 \le x \le a$ and $1 - a \le x \le 1$ and set $f(x) = -L_x h(x)$. Then $f \in C(I)$ also satisfies f(x) = 0 for $0 \le x \le a$ and $1 - a \le x \le 1$.

Since $\mu_n = 1/\lambda_n$ in Theorem 4.3 (Mercer's theorem), the series

$$r(t, x, y) = \sum_{n=1}^{\infty} \frac{e^{-\lambda_n t}}{\lambda_n} \phi_n(x) \phi_n(y)$$
(5.7)

converges uniformly for $0 \le x, y \le 1$ and $0 \le t < \infty$. Also, $r(0, x, y) = \sum_{n=1}^{\infty} (1/\lambda_n)\phi_n(x)\phi_n(y) = g(x, y)$. Since m(dy) = m(y)dy where m(y) is bounded and continuous for $a \le y \le 1 - a$, the function

$$u(x,t) = \int_0^1 r(t,x,y)f(y)m(dy)$$
(5.8)
$$= \int_0^1 \left(\sum_{n=1}^\infty \frac{e^{-\lambda_n t}}{\lambda_n} \phi_n(x)\phi_n(y)\right)f(y)m(dy)$$
$$= \sum_{n=1}^\infty \frac{e^{-\lambda_n t}}{\lambda_n} \phi_n(x)\int_0^1 \phi_n(y)f(y)m(dy)$$

with uniform convergence for $0 \le x \le 1$ and $0 \le t < \infty$. Also

$$u(x,0) = \sum_{n=1}^{\infty} \frac{1}{\lambda_n} \phi_n(x) \int_0^1 \phi_n(y) f(y) m(dy)$$

=
$$\int_0^1 r(0,x,y) f(y) m(dy) = \int_0^1 g(x,y) f(y) m(dy) = h(x)$$

by (4.14) and Theorem 2.1. Thus by (5.7), (4.15), and (4.5), the function u(x,t) in (5.8) is continuous for $0 \le x \le 1$ and $t \ge 0$ and satisfies

$$\frac{\partial}{\partial t}u(x,t) = L_x u(x,t), \qquad 0 < x < 1, \quad 0 < t < \infty$$
(5.9)

$$u(0,t) = u(1,t) = 0, \qquad 0 \le t < \infty$$

$$u(x,0) = h(x), \qquad 0 \le x \le 1$$
(5.10)

It follows from (4.1) and g(x, y) = g(y, x) that

$$\int_{0}^{1} \phi_{n}(x)h(x)m(dx) = \int_{0}^{1} \int_{0}^{1} \phi_{n}(x)g(x,y)m(dx) f(y)m(dy)$$
$$= \frac{1}{\lambda_{n}} \int_{0}^{1} \phi_{n}(y)f(y)m(dy)$$

Hence by (5.8)

$$u(x,t) = \sum_{n=1}^{\infty} e^{-\lambda_n t} \phi_n(x) \int_0^1 \phi_n(y) h(y) m(dy)$$

= $\int_0^1 \sum_{n=1}^{\infty} e^{-\lambda_n t} \phi_n(x) \phi_n(y) h(y) m(dy)$
= $\int_0^1 p(t,x,y) h(y) m(dy)$ (5.11)

By maximum principles for parabolic linear partial differential equations, any solution of (5.9) that is continuous for $0 \le x \le 1$ and $0 \le t \le T$ satisfies

$$\max_{0 \le x \le 1, \ 0 \le t \le T} u(x,t)$$
(5.12)
=
$$\max \left\{ \max_{0 \le t \le T} u(0,t), \ \max_{0 \le t \le T} u(1,t), \ \max_{0 \le x \le 1} u(x,0) \right\}$$

with a similar inequality for the minimum. (See Lemma 3, p167, of Protter and Weinberger, 1967.) If $0 \le h(x) \le 1$ in (5.10), this implies $0 \le u(x,t) \le 1$ for $0 \le x \le 1$ and $0 < t < \infty$.

We have now shown that u(x,t) in (5.11) satisfies (5.12) for any $h \in C^2(I)$ with h(x) = 0 for $0 \le x \le a$ and $1 - a \le x \le 1$ (some a > 0). If either $p(t_0, x_0, y_0) < 0$ or $\int_0^1 p(t_0, x_0, y)m(dy) > 1$ for any $t_0 > 0$ and $0 < x_0, y_0 < 1$, we quickly arrive at contradiction. This completes the proofs of parts (i) and (ii) of Theorem 5.1.

If f(x) is smooth and vanishes near the endpoints, part (iii) of Theorem 5.1 follows from (5.11) and the arguments from (5.8) to (5.10). By part (ii)

$$|Q_t f(x)| \leq \int_0^1 p(t, x, y) |f(y)| m(dy) \leq ||f|| = \max_{0 \leq y \leq 1} |f(y)|$$
(5.13)

for $f \in C[0, 1]$. The uniformity in (5.13) implies that if (5.6) holds for all h is a subset $\mathcal{L} \subset C_0(I_0)$, then it also holds for all $h \in \overline{\mathcal{L}}$ with closure in the uniform norm on (0, 1). Thus (5.6) follows for all $f \in C_0(I_0)$. This completes the proof of Theorem 5.1.

6. Hille-Yoshida Theory. The linear operators Q_t defined in (5.4) satisfy $Q_{t+s} = Q_t Q_s$ by (5.5) and

$$\lim_{t \to 0} \|Q_t f - f\| = 0 \quad \text{for all } f \in C_0(I_0)$$
(6.1)

by (5.6) where $||f|| = \max_{0 \le x \le 1} |f(x)|$. In other works, Q_t is a stronglycontinuous semi-group of linear operators on the Banach space $B = C_0(I_0)$.

The infinitesimal generator of a strongly continuous semigroup of linear operators on a Banach space B is defined to be the linear operator on the subspace $\mathcal{D} = \mathcal{D}(A) \subseteq B$ defined by

$$\mathcal{D}(A) = \{ h \in B : \lim_{t \to 0} \| (1/t)(Q_t h - h) - f \| \text{ for some } f \in B \}$$
(6.2)

with $Ah = \underline{f}$. The operator A is linear on $\mathcal{D}(A)$, but generally $\mathcal{D}(A) \subset B$. In general $\overline{\mathcal{D}(A)} = B$ if Q_t is strongly continuous, since, for any a > 0 and $f \in B$, $h = \int_0^a Q_s f ds \in \mathcal{D}(A)$ with $Ah = Q_a f - f$.

If $||Q_{t_0}f|| \leq C||f||$ for all $f \in B$ and any $t_0 > 0$ and $C < \infty$, then exists constants $C_1 > 0$ and real K such that

$$\|Q_t f\| \le C_1 e^{Kt} \|f\| \tag{6.3}$$

for all $t < \infty$ and $f \in B$. If $\lambda > K$, the resolvent operator

$$R_{\lambda}f = \int_0^\infty e^{-\lambda t} Q_t f dt \tag{6.4}$$

is then a bounded linear operator on B.

Lemma 6.1. Let Q_t be a semigroup of linear operators on a Banach space B satisfying (6.1) and (6.3). Then, for any $\lambda > K$ for K in (6.3), R_{λ} maps B onto $\mathcal{D}(A)$ and

$$\lambda h - Ah = f$$
 whenever $h = R_{\lambda} f$

Proof. If $h = R_{\lambda} f$ for $f \in B$, then

$$Q_t h = \int_0^\infty e^{-\lambda s} Q_{t+s} f ds = e^{\lambda t} \int_t^\infty e^{-\lambda s} Q_s f ds$$
$$= e^{\lambda t} h - e^{\lambda t} \int_0^t e^{-\lambda s} Q_s f ds$$

Thus $(1/h)(Q_th - h) \to \lambda h - f$ as $t \to 0$. This implies that, for any $f \in B$, $h = R_{\lambda}f \in \mathcal{D}(A)$ and $Ah = \lambda h - f$, or $(\lambda - A)h = f$. In particular, (i) the range of R_{λ} is contained in the domain $\mathcal{D}(A)$, (ii) the range of $\lambda I - A$ is all of B, and (iii) $(\lambda - A)R_{\lambda}f = f$ for all $f \in B$.

Conversely, assume $h_1 \in \mathcal{D}(A)$. Set $f = (\lambda - A)h_1$ and $h = R_{\lambda}f$. Then $(\lambda - A)h = f = (\lambda - A)h_1$ and, if $g = h - h_1$, $g \in \mathcal{D}(A)$ and $(\lambda - A)g = 0$. In general if $g \in \mathcal{D}(A)$ and s > 0, then

$$(1/t)(Q_t - I)Q_s g = (1/t)(Q_{t+s} - Q_s)g = Q_s((1/t)Q_t g - g) \to Q_s Ag$$

as $t \to 0$. This implies that $Q_s g \in \mathcal{D}(A)$ and $AQ_s = Q_s A$. If $Ag = \lambda g$, then $Q_t g \in \mathcal{D}(A)$ for all t > 0 and $A(Q_t g) = Q_t Ag = \lambda Q_t g$. This implies $Q_t g = e^{\lambda t} g$, which violates (6.3) if $g \neq 0$ and $\lambda > K$. Thus g = 0, which implies $h_1 = h$. Hence the range of R_λ is all of $\mathcal{D}(A)$, which completes the proof of the lemma.

We can find A and $\mathcal{D}(A)$ exactly for the semigroup Q_t defined by (5.4) in Section 5 for $B = C_0(I_0)$.

Theorem 6.1. Let $L_x = (d/dm(x))(d/ds(x))$ be the operator in (1.1) and let Q_t be as in (5.4) for $B = C_0(I_0)$. Then

$$\mathcal{D}(A) = \{ h : h \in C^2(I_0) \cap C_0(I_0) \text{ and } L_x h \in C_0(I_0) \}$$
(6.6)

with

$$Ah(x) = L_x h(x) = \frac{d}{dm(x)} \frac{d}{ds(x)} h(x)$$

Proof. It follows from (5.1), (4.11), and Theorem 5.1 that

$$|Q_t f(x)| = \left| \int_0^1 p(t, x, y) f(y) m(dy) \right| \le C \sum_{n=1}^\infty \lambda_n^2 e^{-\lambda_n t} \|f\|$$

for $f \in C(I)$ and $0 \leq t < \infty$. Since $\psi(t) = \left(\sum_{n=1}^{\infty} \lambda_n^2 e^{-\lambda_n t}\right)/e^{-\lambda_1 t}$ is bounded and continuous for $1 \leq t < \infty$ with $\lim_{t\to\infty} \psi(t) = C$ where C is positive and finite, it follows that

$$|Q_t f(x)| \leq C e^{-\lambda_1 t} ||f||$$

for $0 \le t < \infty$. Thus (6.3) holds with $K = -\lambda_1$. Hence by Lemma 6.1 with $\lambda = 0$

$$R_0 f(x) = \int_0^\infty Q_s f(s) ds = \int_0^1 g(x, y) f(y) m(dy)$$

by (5.2). Hence $\mathcal{D}(A)$ is the range of R_0 on $B = C_0(I_0)$, which by Theorem 2.1 is exactly $\mathcal{D}(A)$ in (6.6). By Lemma 6.1, $h \in \mathcal{D}(A)$ if and only if $h = R_0 f$ for some $f \in B$ with Ah = -f, so that $Ah = L_x h$ by Theorem 2.1.

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