## VII.C. Inner products

Let $V / \mathbb{R}$ be an $n$-dimensional real vector space. Recall from the end of §VII.B that a symmetric bilinear form

$$
B: V \times V \rightarrow \mathbb{R}
$$

is called positive-definite if it is nondegenerate of signature $(p, q)=$ $(n, 0)$ - or equivalently, if

$$
\vec{v} \neq \overrightarrow{0} \quad \Longrightarrow \quad B(\vec{v}, \vec{v})>0 .
$$

VII.C.1. DEFINITION. An inner product is a positive-definite symmetric bilinear form.

Notation: Instead of $B(\vec{u}, \vec{w})$ we shall write $\langle\vec{u}, \vec{w}\rangle$. Since $\langle\cdot, \cdot\rangle$ is definite,

$$
0=\|\vec{v}\|^{2}:=\langle\vec{v}, \vec{v}\rangle \quad \Longrightarrow \quad \vec{v}=\overrightarrow{0} .
$$

A given inner product determines notions of length and orthogonality on $V$ :

$$
\vec{u} \perp \vec{w} \Leftrightarrow\langle\vec{u}, \vec{w}\rangle=0, \quad\|\vec{u}\|=\sqrt{\langle\vec{u}, \vec{u}\rangle} .
$$

If $V=\mathbb{R}^{n}$ these should not be mixed with the $\|\cdot\|$ and $\perp$ coming from the dot product. ${ }^{11}$

For any symmetric bilinear form $B$, Sylvester's theorem guarantees the existence of a basis $\mathcal{B}$ such that

$$
B(\vec{u}, \vec{w})={ }^{t}[\vec{u}]_{\mathcal{B}}\left(\begin{array}{ccc}
\mathbb{I}_{p} & 0 & 0 \\
0 & -\mathbb{I}_{q} & 0 \\
0 & 0 & 0
\end{array}\right)[\vec{w}]_{\mathcal{B}} .
$$

Clearly in order for $B(\vec{u}, \vec{u})$ to be always positive ( $\vec{u} \neq 0$ ) we must have $q=0$ and $p=n$ (all 1's along the diagonal). Therefore any given inner product has a basis with respect to which

$$
\langle\vec{u}, \vec{w}\rangle={ }^{t}[\vec{u}]_{\mathcal{B}}[\vec{w}]_{\mathcal{B}}=\left\{\begin{array}{c}
\text { "dot product" in that } \\
\text { coordinate system }
\end{array} .\right.
$$

[^0]WARNING: If $V=\mathbb{R}^{n}$ this is not necessarily the usual dot product, unless $\mathcal{B}=\hat{e}$ !

If instead $V$ is a vector space over $\mathbb{C}$, then by "inner product" we shall mean a positive-definite Hermitian form $\langle\cdot, \cdot\rangle: V \times V \rightarrow \mathbb{C}$. In particular, the $1^{\text {st }}$ entry is conjugate-linear $[\langle\alpha \vec{u}, \vec{w}\rangle=\bar{\alpha}\langle\vec{u}, \vec{w}\rangle]$, the $2^{\text {nd }}$ linear $[\langle\vec{u}, \alpha \vec{w}\rangle=\alpha\langle\vec{u}, \vec{w}\rangle]$, and $\langle\vec{u}, \vec{w}\rangle=\overline{\langle\vec{w}, \vec{u}\rangle}$. By Sylvester there is a basis with respect to which

$$
\langle\vec{u}, \vec{w}\rangle=[\vec{u}]_{\mathcal{B}}^{*}[\vec{w}]_{\mathcal{B}},
$$

which is "like" the complex dot product (and the same warning applies when $V=\mathbb{C}^{n}$ ). By "plugging in" the basis elements for $\vec{u}, \vec{w}$ one has:
VII.C.2. Proposition. For any inner product on a complex (real) vector space $V$, there is a basis $\mathcal{B}=\left\{\vec{v}_{1}, \ldots, \vec{v}_{n}\right\}$ for $V$ such that $\left\langle\vec{v}_{i}, \vec{v}_{j}\right\rangle=$ $\delta_{i j}$, i.e. $\mathcal{B}$ is a unitary (orthonormal) basis under the inner product.

WARNING: For $\mathbb{C}^{n}\left(\right.$ resp. $\left.\mathbb{R}^{n}\right), \hat{e}$ is unitary (orthonormal) only under the standard inner $(=\operatorname{dot})$ product.
VII.C.3. Remark. $V$ together with some inner product $\langle\cdot, \cdot\rangle$ (and the accompanying notions of $\perp$ and length) is called an "inner product space". The Proposition simply says: any inner product space has a unitary (orthonormal) basis, though (as the next Example attests) this basis is far from unique. The practical way to obtain one, starting from any given basis for $V$, is by the Gram-Schmidt procedure (exactly as we did it for the dot product); see for example Exercise III.B.2(a).
VII.C.4. Example. Consider the symmetric bilinear form

$$
B(\vec{x}, \vec{y})=\langle\vec{x}, \vec{y}\rangle={ }^{t} \vec{x} M \vec{y}:=\left(\begin{array}{ll}
x_{1} & x_{2}
\end{array}\right)\left(\begin{array}{cc}
\frac{1}{8} & 0 \\
0 & \frac{1}{2}
\end{array}\right)\binom{y_{1}}{y_{2}}=\frac{x_{1} y_{1}}{8}+\frac{x_{2} y_{2}}{2}
$$

on $V=\mathbb{R}^{2}$ (with vectors written with respect to the standard basis). Clearly this is positive definite, hence an inner product, and $\mathcal{B}=$
$\left\{\vec{v}_{1}, \vec{v}_{2}\right\}:=\left\{\binom{\sqrt{8}}{0},\binom{0}{\sqrt{2}}\right\}$ is an orthonormal basis of this inner product space, i.e. $\left\langle\vec{v}_{i}, \vec{v}_{j}\right\rangle=\delta_{i j}$.

For any matrix $S$ satisfying ${ }^{12}$ t $S M S=M$, we have

$$
\left\langle S \vec{v}_{i}, S \vec{v}_{j}\right\rangle={ }^{t} \vec{v}_{i}^{t} S M S \vec{v}_{j}={ }^{t} \vec{v}_{i} M \vec{v}_{j}=\left\langle\vec{v}_{i}, \vec{v}_{j}\right\rangle=\delta_{i j}
$$

hence $S \mathcal{B}=\left\{S \vec{v}_{1}, S \vec{v}_{2}\right\}$ is another orthonormal basis. These matrices are precisely the ones of the form

$$
S=\left(\begin{array}{cc}
\cos \theta & -2 \sin \theta \\
\frac{1}{2} \sin \theta & \cos \theta
\end{array}\right) .
$$

So taking $\theta=-\frac{\pi}{4}$, we get $S=\left(\begin{array}{cc}\sqrt{2} / 2 & \sqrt{2} \\ -\sqrt{2} / 4 & \sqrt{2} / 2\end{array}\right)$ and $S \mathcal{B}=\left\{\binom{2}{-1},\binom{2}{1}\right\}$, which as you can check, is indeed orthonormal in $\langle\cdot, \cdot\rangle$.

Generalizing one more notion from the dot product setting to an arbitrary complex (resp. real) inner product space, we have the
VII.C.5. Definition. Call a transformation $T: V \rightarrow V$ unitary (resp. orthogonal) in the given inner product if $\|T \vec{v}\|=\|\vec{v}\|$ for all $\vec{v} \in V$. (Equivalently, $\langle T \vec{v}, T \vec{w}\rangle=\langle\vec{v}, \vec{w}\rangle \forall \vec{v}, \vec{w} \in V$; see Exercise (5).)

Multiplication by $S$ in Example VII.C. 4 is orthogonal in this sense.
Gram coefficients, Fourier expansion. Next let $(V,\langle\cdot, \cdot\rangle)$ be an inner product space over $\mathbb{C}$ (resp. $\mathbb{R}$ ), $\mathcal{B}=\left\{\vec{v}_{1}, \ldots, \vec{v}_{n}\right\}$ a unitary (orthonormal) basis, and $T: V \rightarrow V$ a linear transformation. Write $a_{i j}$ for the entries of $A=[T]_{\mathcal{B}}$ so that

$$
T \vec{v}_{j}=\sum_{i} a_{i j} \vec{v}_{i},
$$

and notice that

$$
\begin{gathered}
\left\langle\vec{v}_{i}, T \vec{v}_{j}\right\rangle=\left\langle\vec{v}_{i}, \sum_{k} a_{k j} \vec{v}_{k}\right\rangle=\sum_{k} a_{k j}\left\langle\vec{v}_{i}, \vec{v}_{k}\right\rangle=\sum_{k} a_{k j} \delta_{i k}=a_{i j} \\
\Longrightarrow T \vec{v}_{j}=\sum_{i}\left\langle\vec{v}_{i}, T \vec{v}_{j}\right\rangle \vec{v}_{i} \Longrightarrow \quad \Longrightarrow \quad T \vec{x}=\sum_{i}\left\langle\vec{v}_{i}, T \vec{x}\right\rangle \vec{v}_{i},
\end{gathered}
$$

where the last implication is by linear extension. Notice that (for $T=$ Id) we've recovered the Fourier expansion formula $\vec{x}=\sum_{i}\left\langle\vec{v}_{i}, \vec{x}\right\rangle \vec{v}_{i}$
${ }^{12}$ That is, in the notation of §VII.B, $S \in \mathrm{O}_{2}(\mathbb{R}, B)$; we have $\langle S \vec{x}, S \vec{y}\rangle=\langle\vec{x}, \vec{y}\rangle$, and $S$ preserves the inner product/bilinear form.
from §VII.A, but for a general inner product. The coefficients $\left\langle\vec{v}_{i}, T \vec{v}_{j}\right\rangle$ are called Gram coefficients. WARNING: They only compute the entries of $[T]_{\mathcal{B}}$ if $\mathcal{B}$ is unitary (orthonormal).

Adjoint of a linear transformation. Now we define a new transformation, ${ }^{13}$ the adjoint

$$
T^{\dagger}: V \rightarrow V
$$

of $T$, by demanding that for all $\vec{u}, \vec{w}$

$$
\begin{equation*}
\left\langle T^{\dagger} \vec{u}, \vec{w}\right\rangle=\langle\vec{u}, T \vec{w}\rangle . \tag{VII.C.6}
\end{equation*}
$$

How do we know such a transformation exists? If it did then we could apply the expansion formula to obtain ${ }^{14}$

$$
T^{+} \vec{x}=\sum_{i}\left\langle\vec{v}_{i}, T^{+} \vec{x}\right\rangle \vec{v}_{i}=\sum_{i} \overline{\left\langle T^{+} \vec{x}, \vec{v}_{i}\right\rangle} \vec{v}_{i}=\sum_{i} \overline{\left\langle\vec{x}, T \vec{v}_{i}\right\rangle} \vec{v}_{i}
$$

and the latter expression is as good as a definition (plug it in to (VII.C.6) and check). Since

$$
\left.T^{\dagger} \vec{v}_{j}=\sum \overline{\left\langle\vec{v}_{j}, T \vec{v}_{i}\right\rangle}\right\rangle \vec{v}_{i}=\sum \bar{a}_{j i} \vec{v}_{i}
$$

the matrix $\left[T^{\dagger}\right]_{\mathcal{B}}$ is the conjugate transpose of $[T]_{\mathcal{B}}$ (or just the transpose if $V / \mathbb{R}$ ). WARNING: While this matrix point of view on the adjoint is very important, you can't just write the matrix of $T$ with respect to an arbitrary basis, conjugate transpose it and claim you found the matrix of $T^{\dagger}$. (If that were possible then $T^{\dagger}$ would have nothing to do with the inner product.) As usual, $\left[T^{\dagger}\right]_{\mathcal{B}}={ }^{t} \overline{[T]_{\mathcal{B}}}$ only holds for $\mathcal{B}$ unitary/orthogonal (under the given inner product).

Here's a slightly more abstract point of view on what the adjoint "is": for fixed $\vec{w} \in V,\langle\vec{w}, \cdot\rangle$ is a linear functional on $V$ (defined by $\langle\vec{w}, \cdot\rangle(\vec{v})=\langle\vec{w}, \vec{v}\rangle)$, i.e. an element of $V^{\vee} .{ }^{15}$ In fact the map sending $\vec{w} \mapsto\langle\vec{w}, \cdot\rangle$ gives an isomorphism of $V$ with its dual space.
$\overline{13}$ not to be confused with the classical adjoint (or adjugate) $\operatorname{Adj}(M)$ of a matrix $M$.
${ }^{14}$ This computation and those that follow are valid for $V$ over either $\mathbb{R}$ or $\mathbb{C}$. You may ignore the bars (denoting complex conjugation) in the real case, since the numbers are all real.
${ }^{15}$ Recall from §III.D that the symbol " $\vee$ " dualizes vector spaces and linear transformations.

QUESTION: What transformation on $V$ corresponds under this isomorphism to the dual transformation $T^{\vee}: V^{\vee} \rightarrow V^{\vee}$ ? That is, what dotted map makes the following diagram "commute"?


The answer: $T^{\dagger}$. In other words, the top map followed by (composed with) $T^{\vee}$ is the same as $T^{\dagger}$ followed by the bottom map:

$$
\left\langle T^{\dagger} \vec{w}, \cdot\right\rangle=T^{\vee}\langle\vec{w}, \cdot\rangle
$$

as elements of $V^{\vee}$. One sees this by applying both sides to $\vec{v} \in V$ : using the definition of $T^{\vee}$,

$$
\left(T^{\vee}\langle\vec{w}, \cdot\rangle\right)(\vec{v})=\langle\vec{w}, \cdot\rangle(T \vec{v})=\langle\vec{w}, T \vec{v}\rangle=\left\langle T^{\dagger} \vec{w}, \vec{v}\right\rangle=\left\langle T^{\dagger} \vec{w}, \cdot\right\rangle(\vec{v}) .
$$

CONCLUSION: Let $V$ be an inner product space and $T: V \rightarrow V a$ linear transformation. Then if we use the inner product $\langle\cdot, \cdot\rangle: V \times V \rightarrow \mathbb{C}$ to view $V$ as its own dual space, $T^{\dagger}: V \rightarrow V$ plays the role of $T^{\vee}$.

Since $\left(T^{\vee}\right)^{\vee}=T$ you might think $\left(T^{\dagger}\right)^{\dagger}=T$, and this is correct: write for any $\vec{w}, \vec{u} \in V$

$$
\langle\vec{w}, T \vec{u}\rangle=\left\langle T^{\dagger} \vec{w}, \vec{u}\right\rangle=\overline{\left\langle\vec{u}, T^{\dagger} \vec{w}\right\rangle}=\overline{\left\langle\left(T^{\dagger}\right)^{\dagger} \vec{u}, \vec{w}\right\rangle}=\left\langle\vec{w},\left(T^{\dagger}\right)^{\dagger} \vec{u}\right\rangle ;
$$

thus

$$
\left\langle\vec{w}, T \vec{u}-\left(T^{\dagger}\right)^{\dagger} \vec{u}\right\rangle=0 \quad \forall \vec{u}, \vec{w} .
$$

In particular, by taking $\vec{w}=T \vec{u}-\left(T^{\dagger}\right)^{\dagger} \vec{u}$ we have $\left\|T \vec{u}-\left(T^{\dagger}\right)^{\dagger} \vec{u}\right\|=$ 0 , and thus $T \vec{u}-\left(T^{\dagger}\right)^{\dagger}=0$ or $T \vec{u}=\left(T^{\dagger}\right)^{\dagger} \vec{u}$ for all $\vec{u} \in V$.
VII.C.7. Remark. The adjoint gives another way to characterize unitary transformations (Defn. VII.C.5). From $\langle T \vec{x}, T \vec{y}\rangle=\langle\vec{x}, \vec{y}\rangle$ $(\forall \vec{x}, \vec{y} \in V)$ we get $\left\langle T^{\dagger} T \vec{x}, \vec{y}\right\rangle=\langle\vec{x}, \vec{y}\rangle$ hence $\left\langle\left(T^{\dagger} T-\mathbb{I}\right) \vec{x}, \vec{y}\right\rangle=0$. Taking $\vec{y}=\left(T^{\dagger} T-\mathbb{I}\right) \vec{x}$ gives $\left\|\left(T^{\dagger} T-\mathbb{I}\right) \vec{x}\right\|=0$ hence $\left(T^{\dagger} T-\mathbb{I}\right) \vec{x}=\overrightarrow{0}$ for all $\vec{x}$. So $T^{\dagger} T-\mathbb{I}=0$ whence

$$
\begin{equation*}
T^{\dagger} T=\mathbb{I}=T T^{\dagger} \tag{VII.C.8}
\end{equation*}
$$

VII.C.9. EXAMPLE. Let $V$ be the real vector space of all continuously differentiable (real-valued) functions on $\mathbb{R}$ which are periodic of period 1: i.e. $f(t)=f(t+1)$ for all $t$. (This is infinite-dimensional, but no matter.) Then ${ }^{16}$

$$
\langle f, g\rangle=\int_{0}^{1} f(t) g(t) d t
$$

defines an inner product: it is clearly bilinear, and for any nonzero function $f$ we have $\langle f, f\rangle=\int_{0}^{1} f(t)^{2} d t>0$. Now define a transformation $T: V \rightarrow V$ by $T(f)=3 f-f^{\prime}$ (i.e., $T=3-\frac{d}{d t}$ ). We use integration by parts to find the adjoint:

$$
\begin{gathered}
\langle f, T g\rangle=\int_{0}^{1} f\left(3 g-g^{\prime}\right) d t=3 \int_{0}^{1} f g d t-\int_{0}^{1} f g^{\prime} d t \\
=3 \int_{0}^{1} f g d t+\int_{0}^{1} f^{\prime} g d t=\int_{0}^{1}\left(3 f+f^{\prime}\right) g d t=\left\langle T^{\dagger} f, g\right\rangle
\end{gathered}
$$

that is, $T^{\dagger}=3+\frac{d}{d t}$. (There was no $f(1) g(1)-f(0) g(0)$ term in the $\int$-by-parts because this is zero by the periodicity hypothesis!)

## Exercises

(1) Let $V=\mathbb{R}^{2}$, and specify an inner product $\langle\cdot, \cdot\rangle$ by

$$
\|\vec{x}\|^{2}:=\left(x_{1}-x_{2}\right)^{2}+3 x_{2}^{2}
$$

(Here $\vec{x}=x_{1} \hat{e}_{1}+x_{2} \hat{e}_{2}$.) Find the matrix, with respect to $\hat{e}$, of the orthogonal projection (orthogonal with respect to $\langle\cdot, \cdot\rangle$, not the dot product) onto the line generated by $3 \hat{e}_{1}+4 \hat{e}_{2}$. Show that this projection is its own adjoint.
(2) Check that $T^{\dagger} \vec{x}:=\sum_{i} \overline{\left\langle\vec{x}, T \vec{v}_{i}\right\rangle} \vec{v}_{i}$ satisfies equation (VII.C.6) as claimed in the notes. (Here $\left\{\vec{v}_{i}\right\}_{i=1}^{n}$ is a unitary/orthonormal basis.)
(3) In Exercise VII.B.2, you found an orthonormal basis $\mathcal{B}$ of $\mathcal{P}_{2}(\mathbb{R})$ under $\langle f, g\rangle:=\int_{0}^{1} f(t) g(t) d t$, and computed the matrix $[T]_{\mathcal{B}}$, where $(T f)(t):=f(t-1)$.

[^1](a) Find $\left[T^{\dagger}\right]_{\mathcal{B}}$, and compute $T^{\dagger}(1)$. Check your answer by "plugging it in" to the definition for adjoint. It should be a quadratic polynomial with rather nasty coefficients.
(b) What is the adjoint of $S:=\frac{d}{d t}$ ?
(4) (a) For $V=M_{n}(\mathbb{C})$, check that $\langle A, B\rangle:=\operatorname{tr}\left(A B^{*}\right)$ defines an inner product.
(b) Now let $P$ be a fixed invertible matrix in $V$, and define $T_{P}$ : $V \rightarrow V$ by $T_{P}(A):=P^{-1} A P$. What is $T_{P}^{+}$?
(5) By definition, unitary / orthogonal transformations preserve $\|\cdot\|$; show that this implies they preserve $\langle\cdot, \cdot\rangle$. [Hint: in the orthogonal (real) case you can write $\langle\vec{x}, \vec{y}\rangle=\frac{1}{2}\langle\vec{x}+\vec{y}, \vec{x}+\vec{y}\rangle-\frac{1}{2}\langle\vec{x}, \vec{x}\rangle-$ $\frac{1}{2}\langle\vec{y}, \vec{y}\rangle$. What do you have to do differently in the unitary (complex) case?]


[^0]:    $\overline{11}$ often called the "Euclidean" or "standard" inner product on $\mathbb{R}^{n}$.

[^1]:    ${ }^{16}$ If we considered complex-valued functions, then $\langle f, g\rangle$ would be $\int_{0}^{1} \overline{f(t)} g(t) d t$.

