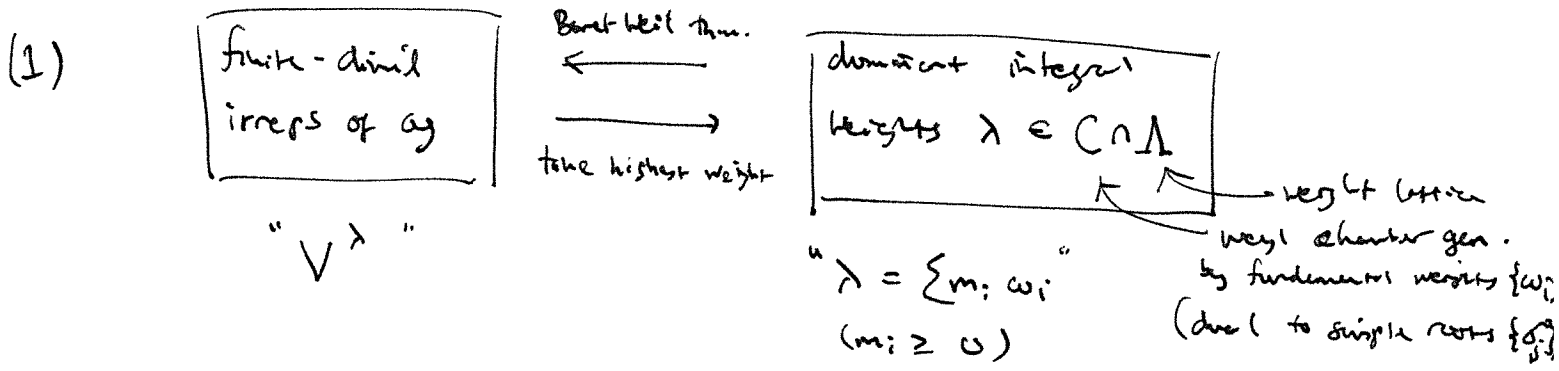


6. Homogeneous VHS and cycles on abelian varieties

We now give a second application of infinitesimal invariants.

To begin, we construct the homogeneous variations of Hodge structure (over locally symmetric varieties). Let $\mathfrak{g}_{\mathbb{R}}$ be a semisimple Lie algebra with complexification \mathfrak{g} , and compact Cartan subalgebra $\mathfrak{t}_{\mathbb{R}}$; write $\Delta = \Delta_c \sqcup \Delta_n$ for the decomposition into compact and noncompact roots, and pick $E \in \mathfrak{t}_{\mathbb{R}}$ with $\frac{1}{2}E(\Delta_c) \subset 2\mathbb{Z}$, $\frac{1}{2}E(\Delta_n) \subset 2\mathbb{Z} + 1$.

The Theorem of the highest weight gives a 1-to-1 correspondence



So the (finite-dim!) irreps. of $\mathfrak{g}_{\mathbb{R}}$ take the form $(d\rho_\lambda, \tilde{V}^\lambda)$,

(2)

$$\tilde{V}_{\mathbb{C}}^\lambda = \begin{cases} V^\lambda & (\text{"real" case}) \\ V^\lambda \oplus V^{\tau(\lambda)} & \left\{ \begin{array}{l} \tau(\lambda) = \lambda \text{ ("quaternionic" case)} \\ \tau(\lambda) \neq \lambda \text{ ("complex" case)} \end{array} \right. \\ \left(\frac{\parallel}{\sqrt{2}} \right) \end{cases}$$

where $\tau := -w_0$ ($w_0 \in W = W(G_{\mathbb{C}}, \mathbb{E})$ the longest element). Assuming $E(\lambda) \in 2\mathbb{Z} + 1$, the decomposition

(3)

$$\tilde{V}_{\mathbb{C}}^\lambda = \bigoplus_{p \in \mathbb{Z}} E_{(2p+1)}(d\rho_\lambda(E)) =: \bigoplus_{p \in \mathbb{Z}} \left(\tilde{V}_{\mathbb{C}}^\lambda \right)^{p, -p-1}$$

$\leftarrow \left\{ \text{eigenspace w/ eigenvalue } 2p+1 \text{ for } d\rho_\lambda(E) \right\}$

defines an \mathbb{R} -HS of weight (-1) and level $-E(\lambda)$ on \tilde{V}^λ .

(One can show \tilde{V}^λ is polarized, by a unique \mathfrak{g} -invariant alternating form Q .)

Now take $G =$ semisimple \mathbb{Q} -algebraic group of Hermitian type, s.t. $G_{\mathbb{R}}$ contains a compact Cartan subgroup $T_{\mathbb{R}}$. Choose a cocharacter $\chi_0: G_{m, \mathbb{C}} \rightarrow T_{\mathbb{C}}$ s.t. $E := \chi_0'(1)$ satisfies

$$(4) \quad E(\Delta_c) = 0, \quad E(\Delta_n) = \{\pm 2\}.$$

That is, the $(\text{ad } E)$ -Hodge decomp. on $\mathfrak{g}_{\mathbb{C}}$ takes the form

$$(5) \quad \mathfrak{g}_{\mathbb{C}} = \mathfrak{g}^{-1,1} \oplus \mathfrak{g}^{0,0} \oplus \mathfrak{g}^{1,-1}.$$

Then $\Delta_n \cap \Sigma = \{\delta_I\}$ is a special simple root, i.e. $\lambda_{\text{ad}} = \delta_I + \sum_{j \neq I} m_j \sigma_j$ and $E(\sigma_j) = -2\delta_{Ij}$. [In this way, the choice of I (for amongst

the special nodes of the Dynkin diagram of $\mathfrak{g}_{\mathbb{C}}$) determines the real form $G_{\mathbb{R}}$ of $G_{\mathbb{C}}$.] The $\rho_\lambda \circ \chi_0$ resp. $\text{Ad} \circ \chi_0$ eigenspaces recover the

(compatible) Hodge decompositions (3) resp. (5). To vary them,

compose

$$(6) \quad G_{m, \mathbb{C}} \xrightarrow{\chi_0} T_{\mathbb{C}} \xrightarrow{=: \varphi_0} G_{\mathbb{C}}$$

and take the orbit of (6) under conjugation

$$(7) \quad \mathcal{D} := G(\mathbb{R}) \cdot \varphi_0 \cong G(\mathbb{R}) / \underbrace{G^0(\mathbb{R})}_{\left(\begin{array}{l} \text{centralizer of } \varphi_0; \\ \text{Lie } G^0 = \mathfrak{g}^{0,0} =: \mathfrak{g}^0. \end{array} \right)}$$

This is a Hermitian symmetric domain of $\dim_{\mathbb{C}} \mathcal{D} = \dim \mathfrak{g}^{-1,1}$. Taking

$\Gamma \subseteq G(\mathbb{Z})$ torsion-free of finite index, the quotient

$$(8) \quad \Sigma := \Gamma \backslash \mathcal{D}$$

is a quasi-projective (locally symmetric) variety by the Baily-Borel theorem.

(and assuming G_0 is simple)

Varying the choice of root system and special node, we get the classification of irreducible Hermitian symmetric domains into types

$$(9) \left\{ \begin{array}{l} I_{p,q} \text{ (unitary)}, II_n \text{ (quaternionic)}, III_n (= \text{Siegel } h_n \text{ symmetric}), IV_n \text{ (orthogonal "K3 type")}, \\ EIII \text{ \& } EVII \text{ (exceptional)}. \end{array} \right.$$

Now fix Σ in (8) and $\varphi_0 \in D$ in (6), as well as a

\mathbb{Q} -linear representation $\rho: G \rightarrow \text{Aut}(V, \mathbb{Q})$ s.t. $\rho \circ \varphi_0$ is a HS on V

polarized by \mathbb{Q} . Then the $\{\rho \circ g \varphi_0 g^{-1}\}_{g \in G(\mathbb{R})}$ give a variation of HS

over Σ with geometric monodromy (and derived MTG) G . * Putting this

construction of Hermitian (homogeneous) VHS together with (1)-(3), we

get a bijection

$$(10) \quad \boxed{\begin{array}{l} \text{irreducible} \\ \text{Hermitian (R-VHS)/}\Sigma \\ \text{of weight } (-1) \end{array}} \longleftrightarrow \boxed{\begin{array}{l} \{ \text{dominant integral } \lambda \\ \text{with } E(\lambda) \text{ odd} \} \\ \langle \tau \rangle \end{array}}.$$

" $\tilde{V}_{\mathbb{R}}^{\lambda}$ " " $\lambda = \sum m_i \omega_i$ "

($m_i \geq 0, \sum m_i E(\omega_i) \in 2\mathbb{Z} + 1$)

In all the situations we shall consider, $\tilde{V}_{\mathbb{R}}^{\lambda}$ has an underlying \mathbb{Q} -PVHS \tilde{V}^{λ} .

Recalling the $\{h^k(j)\}$ from §4, which we shall denote

$h_{\lambda}^k(j)$ for the irreducible \mathbb{Q} -VHS $V_{\mathbb{Q}}^{\lambda}$ * (so that in the

"non-real" cases $h_{\tilde{V}^{\lambda}}^k(j) = h_{\lambda}^k(j) \oplus h_{\langle \tau \rangle}^k(j)$), the key point

is that we can compare these via representation theory:

* this looks strange, but the $F^i \subset V_{\mathbb{C}}$ and \mathbb{C} -local system $W_{\mathbb{C}}$ are all \mathbb{F} -conjugate

Theorem 1 (Kostant-K): $H_\lambda^k(j) |_{\mathfrak{g}_0} \cong \bigoplus_{w \in W^0(k,j)} V_0^{w \cdot \lambda}$,

where $|_{\mathfrak{g}_0}$ means the fiber over $\mathfrak{g}_0 \in D$

$$W^0(k,j) = \left\{ w \in W \mid \begin{array}{l} w(\Delta^+) \supseteq \Delta_0^+, |w|=k, \\ \text{and } E(w \cdot \lambda) = 2j+1 \end{array} \right\}$$

$$w \cdot \lambda = w(\lambda + \rho) - \rho, \quad \rho = \frac{1}{2} \sum_{\alpha \in \Delta^+} \alpha = \sum \omega_i$$

V_0^ξ is the irrep. of \mathfrak{g}_0 of highest weight ξ .

(The proof, which I won't do, uses Kostant's theorem on Lie algebra cohomology and the identification $\bigoplus_j H_\lambda^k(j) \cong H^k(\mathfrak{g}_0^{-1}, V^\lambda)$.)
For any λ as in (10) we have the (*)

Corollary: $H_\lambda^0(j) = 0 \quad (\forall j \geq 0)$

$$H_\lambda^1(j) \neq 0 \iff j = \mu(\lambda) := \frac{1}{2}(E(w_I \cdot \lambda) - 1),$$

where w_I is the reflection in the unique simple root σ_I with $E(\sigma_I) \neq 0 (= -2)$.

Therefore, $NF_j(j^* \tilde{V}^\lambda)_\mathbb{Q} = \{0\}$ for any étale neighborhood $\mathcal{T} \rightarrow \mathcal{X}$ if $\mu(\lambda), \mu(\tau(\lambda)) < 0$.

This turns out to rule out normal functions outside of a short list of cases, including (III_3, ω_3) for the Ceresa cycle (of course), part of which is given in Theorem 2 below.

There is one other important point: recall that for $g=3$, the Ceresa cycle was defined on a 2:1 cover of $\mathcal{A}_3 = \{\text{moduli space of 3-dim'l PPAVs}\} \cong \mathfrak{S}_6(\mathbb{Z}) \sqrt[III]{3}$, but does it yet try to

push it down to A_3 . There is a reason for this:

if G has \mathbb{Q} -rank > 1 (i.e. isn't SL_2 or $U(n,1)$, for our purposes), then for any arithmetic $\Gamma \leq G(\mathbb{Q})$ we have

$$(11) \quad 0 = H^1(\Gamma, \tilde{V}^\lambda) = H^1(\Sigma, \tilde{V}^\lambda)$$

by a theorem of Raghunathan. Since admissible normal functions extend to the Zariski closure inside the domain where a VHS is defined (nonsingular), $ANF_U(\tilde{V}^\lambda) = ANF_\Sigma(\tilde{V}^\lambda)$ if Γ is torsion-free, so the topological invariant gives

$$(12) \quad ANF_U(\tilde{V}^\lambda) \hookrightarrow H^1(\Sigma, \tilde{V}^\lambda) \quad \text{for any } U \subset \Sigma_{\text{reg}}$$

Hence $ANF_U(\tilde{V}^\lambda)$ vanishes, and we are forced to look for our normal functions on étale neighborhoods \mathcal{I} (finite covers of U).

Theorem 2 (Keast-K): For D of "tube type", the only pairs (D, λ) s.t. $\tilde{V}^\lambda \rightarrow \Gamma \backslash D$ has the possibility of admitting a nontrivial normal function over an étale neighborhood, are

$$(13) \quad \left\{ \begin{array}{l} (I_{2,2}, \{\omega_3\} + a\omega_2), (I_{3,3}, \omega_3), (II_4, \omega_1 + a\{\omega_4\}), (II_6, \omega_6), \\ (III_1, a\omega_1), (III_2, \omega_1 + a\omega_2), (III_3, \omega_3), (IV_{2n-1}, a\omega_1 + \omega_n), (IV_{2n-2}, a\omega_1 + \{\omega_n\}), \\ \text{and } (E.VII, \omega_7). \end{array} \right.$$

We were able, so far, to construct one new normal function for a case on this list: $(I_{3,3}, \omega_3)$, which is the irreducible VHS of type $(1,9,9,1)$ contained in H^3 of the

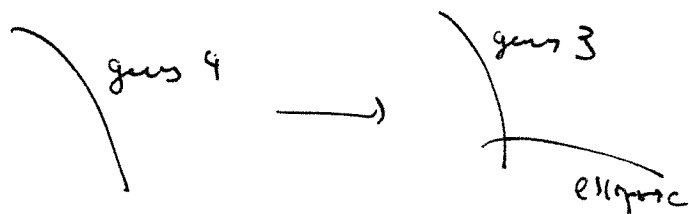
universal Weil abelian 6-fold over $\mathbb{P}^1 \setminus \text{SU}(3,3)/K = \mathcal{B}_{3,3}$. 6

As in the setting of Schoen's proof of the Hodge conjecture for certain Weil 4-folds, a moduli space of generalized Prym varieties dominates $\mathcal{B}_{3,3}$. That is, the general Weil 6-fold parametrized by $\mathcal{B}_{3,3}$ [⊗] arises as the quotient

$$(14) \quad A = J(\tilde{C}) / J(C)$$

where $\tilde{C} \rightarrow C$ is an unramified 3:1 cover. Now
 (genus 10) (genus 4)

one defines the Prym-Ceresa cycle $Z_{\tilde{C}/C} \in \text{Griff}_1(A)_{\mathbb{Q}}$ to be the push-forward of the Ceresa cycle $Z_C \in \text{Griff}_1(J(\tilde{C}))_{\mathbb{Q}}$ to A . By degenerating C



as in Ceresa, we reduce normality of $\overline{AJ^5}(Z_{\tilde{C}/C})$ to that of the Ceresa normal function, thereby producing the desired normal function over a finite cover of $\mathcal{B}_{3,3}$.

⊗ These are Weil 6-folds with multiplication by $\mathbb{Q}(\mathcal{S}_3)$ and a certain polarization invariant