

# Lecture 40 : Change of Variable

Though I tend not to write it out explicitly, we have been doing a lot of substitutions in the integrals we compute. Recall how this goes: given

$$u \longmapsto g(u) = x \quad C^1 \text{ function}$$

mapping

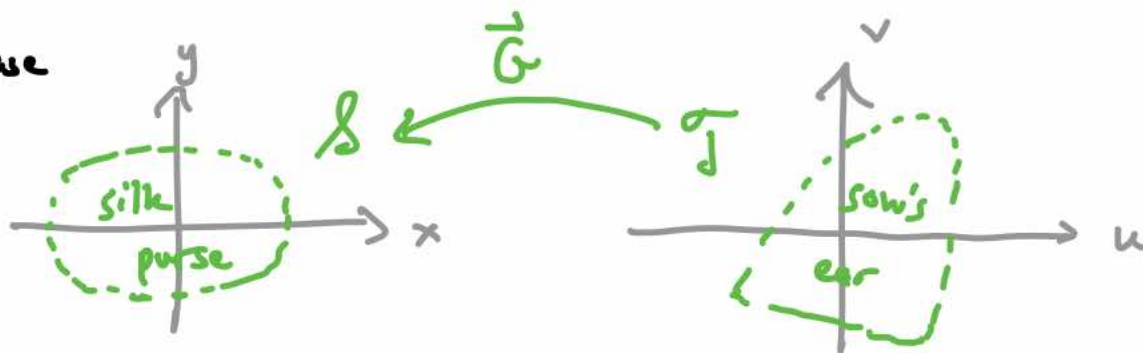
$$[c, d] \longrightarrow [a, b] \quad a = g(c), b = g(d)$$

in 1-to-1 fashion,

$$\int_a^b f(x) dx = \int_c^d f(g(u)) g'(u) du.$$

We are after a 2-variable analogue of this.

Suppose



is a 1-to-1 & onto  $C^1$  map. We'll write

$$(u, v) \longmapsto \vec{G}(u, v) = (x(u, v), y(u, v)),$$

and not worry about  $\vec{G}$  failing to be 1-to-1 on the boundary

or more generally on a content-zero subset. Then for any continuous function  $f: \mathcal{S} \rightarrow \mathbb{R}$ ,

$$\iint_{\mathcal{S}} f \, dA = \iint_{\mathcal{T}} (f \circ \vec{G}) |\det(J_{\vec{G}})| \, dA \quad (*)$$

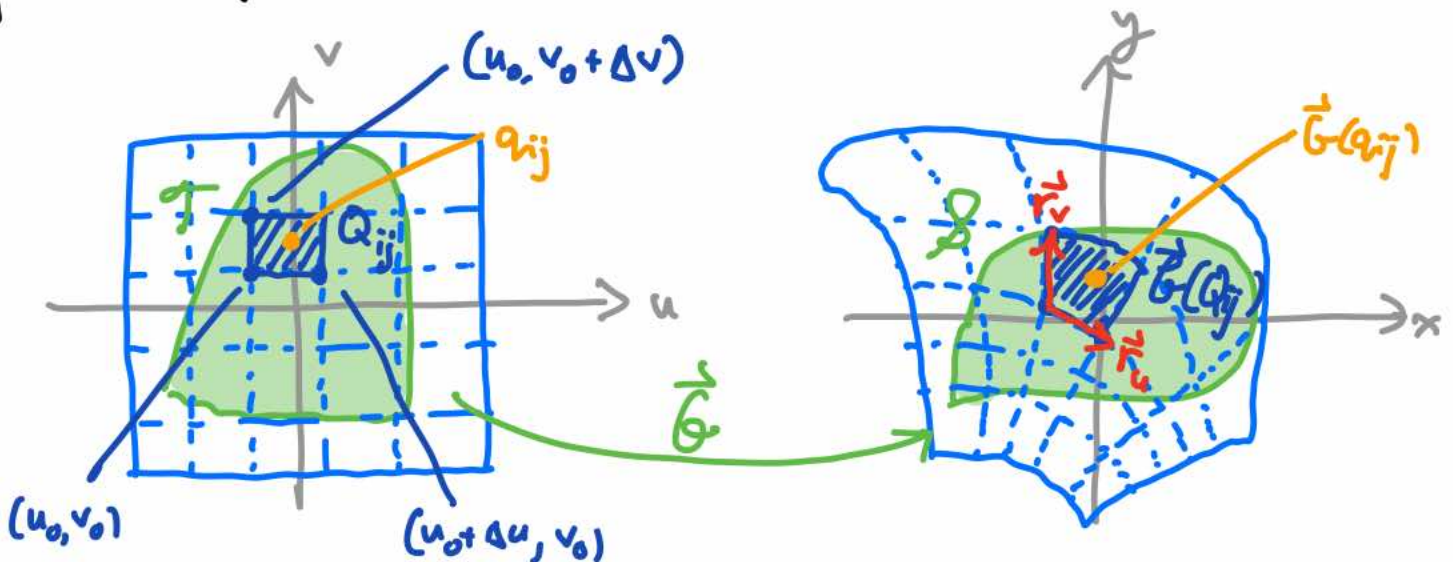
$$\rightarrow \iint_{\mathcal{S}} f(x,y) \, dx \, dy = \iint_{\mathcal{T}} f(\vec{G}(u,v)) |\det(J_{\vec{G}}(u,v))| \, du \, dv$$

where

$$J_{\vec{G}} = \begin{pmatrix} \frac{\partial x}{\partial u} & \frac{\partial x}{\partial v} \\ \frac{\partial y}{\partial u} & \frac{\partial y}{\partial v} \end{pmatrix} \left[ \begin{array}{l} \text{sometimes also written} \\ \frac{\partial(x,y)}{\partial(u,v)} \end{array} \right]$$

is the Jacobian or total derivative of  $\vec{G}$ . We now give an informal explanation of how (\*) is obtained.

Begin by enclosing  $\mathcal{T}$  in a rectangle  $Q$ , and partition  $Q$ :



For simplicity, we assume  $\vec{G}$  extends to  $\mathbb{Q}$ .

The  $\vec{G}(Q_{ij})$  aren't rectangles, but they aren't far from being parallelograms. So we can approximate their area by a cross product:

$$\vec{r}_u := \vec{G}(u_0 + \Delta u, v_0) - \vec{G}(u_0, v_0) \approx J_{\vec{G}}(u_0, v_0) \begin{pmatrix} \Delta u \\ 0 \end{pmatrix} = \begin{pmatrix} \frac{\partial x}{\partial u}(u_0, v_0) \\ \frac{\partial y}{\partial u}(u_0, v_0) \end{pmatrix} \Delta u$$

$$\vec{r}_v := \vec{G}(u_0, v_0 + \Delta v) - \vec{G}(u_0, v_0) \approx J_{\vec{G}}(u_0, v_0) \begin{pmatrix} 0 \\ \Delta v \end{pmatrix} = \begin{pmatrix} \frac{\partial x}{\partial v}(u_0, v_0) \\ \frac{\partial y}{\partial v}(u_0, v_0) \end{pmatrix} \Delta v$$

$$a(\vec{G}(Q_{ij})) \approx \|\vec{r}_u \times \vec{r}_v\| = \left\| \begin{vmatrix} \hat{e}_1 & \hat{e}_2 & \hat{e}_3 \\ \frac{\partial x}{\partial u} \Delta u & \frac{\partial y}{\partial u} \Delta u & 0 \\ \frac{\partial x}{\partial v} \Delta v & \frac{\partial y}{\partial v} \Delta v & 0 \end{vmatrix} \right\|$$

$$= \left| \det \begin{pmatrix} \frac{\partial x}{\partial u} & \frac{\partial y}{\partial u} \\ \frac{\partial x}{\partial v} & \frac{\partial y}{\partial v} \end{pmatrix} \right| \Delta u \Delta v = |\det(J_{\vec{G}}(u_0, v_0))| a(Q_{ij}).$$

So  $|J_{\vec{G}}|$  evaluated at a point of  $Q_{ij}$  approximates the

dilation factor  $\frac{a(\vec{G}(Q_{ij}))}{a(Q_{ij})}$ , and

$$\iint_{\mathcal{D}} f \, dA \stackrel{\text{not rigorously justified}}{=} \lim_{\|P\| \rightarrow 0} \sum_{i,j} f(\vec{G}(q_{ij})) a(\vec{G}(Q_{ij}))$$

*sample point in  $Q_{ij}$*

as the "norm of the partition", i.e. maximum side-length of a  $Q_{ij}$ , goes to 0.

$$= \lim_{\|P\| \rightarrow 0} \sum_{i,j} (f \circ \vec{G})(q_{ij}) |\det(\vec{J}_{\vec{G}}(q_{ij}))| a(Q_{ij})$$

$$= \iint_{\mathcal{J}} (f \circ \vec{G}) |\det(\vec{J}_{\vec{G}})| dA.$$

A key special case on which we shall now focus is

where  $(u, v)$  is called  $(r, \theta)$  and

$$\vec{G}(r, \theta) = (r \cos \theta, r \sin \theta).$$

The Jacobian determinant is

$$|\det(\vec{J}_{\vec{G}})| = \left| \det \begin{pmatrix} \frac{\partial}{\partial r}(r \cos \theta) & \frac{\partial}{\partial \theta}(r \cos \theta) \\ \frac{\partial}{\partial r}(r \sin \theta) & \frac{\partial}{\partial \theta}(r \sin \theta) \end{pmatrix} \right| = \left| \det \begin{pmatrix} \cos \theta & -r \sin \theta \\ \sin \theta & r \cos \theta \end{pmatrix} \right|$$

$$= r \cos \theta (\cos \theta) - (-r \sin \theta) \sin \theta$$

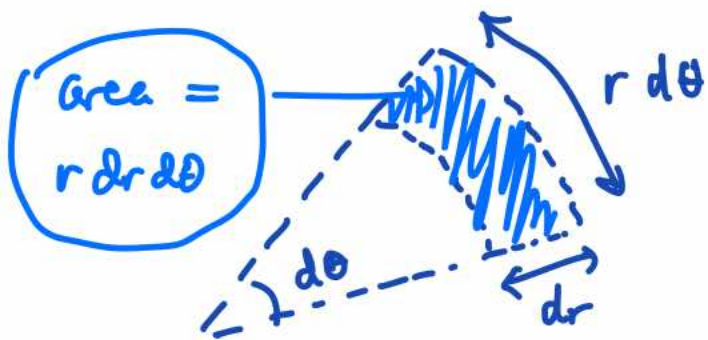
$$= r (\cos^2 \theta + \sin^2 \theta) = r,$$

and so the change-of-variable formula is

$$\boxed{\iint_{\mathcal{L}} f(x, y) dx dy = \iint_{\mathcal{J}} (f \circ \vec{G})(r, \theta) \underline{r} dr d\theta}$$

Intuitively, this is replacing the Riemann sum

over rectangles by one over angular sectors —



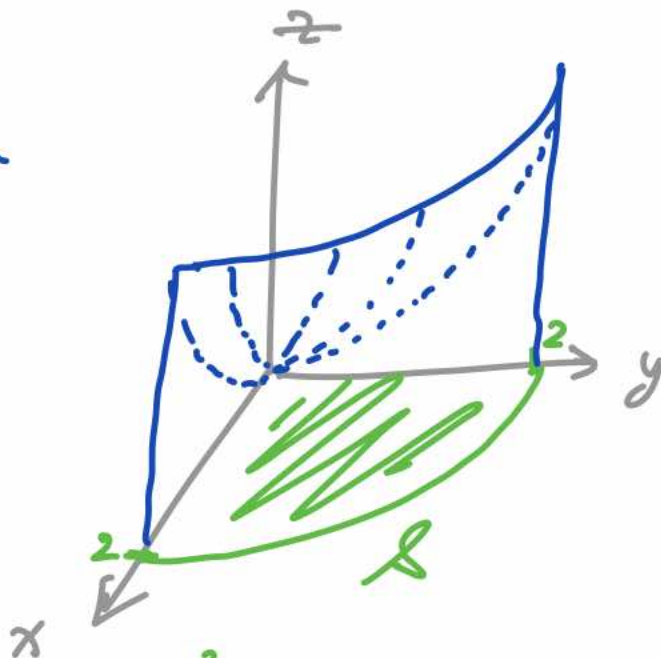
which, after all, are the  $\{G(\Omega_i)\}$  in this case.

We now turn to several

examples of polar integration.

Ex 1 / Find the volume of the solid in the 1<sup>st</sup> octant

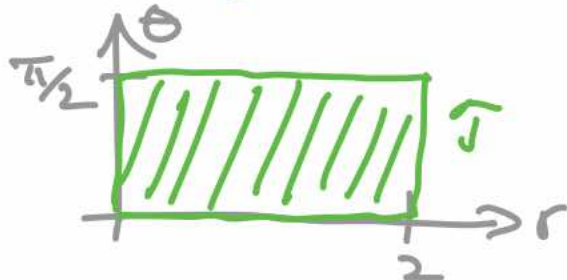
bounded by  $z = x^2 + y^2$ ,  $x^2 + y^2 = 4$ , and the coordinate planes.



$$V = \iint_{\mathcal{L}} (x^2 + y^2) dA = \iint_{\mathcal{L}} \left( (r \cos \theta)^2 + (r \sin \theta)^2 \right) r dr d\theta$$

← what is this?

$$= \int_0^{\pi/2} \left( \int_0^2 r^3 dr \right) d\theta$$



$$= \int_0^{\pi/2} \frac{2^4}{4} d\theta = 4 \cdot \frac{\pi}{2} = 2\pi.$$





Ex 3 / Evaluate  $\int_0^2 \int_0^{\sqrt{2x-x^2}} \sqrt{x^2+y^2} dy dx$ .

This is nasty to antidifferentiate in  $x$  or  $y$ .

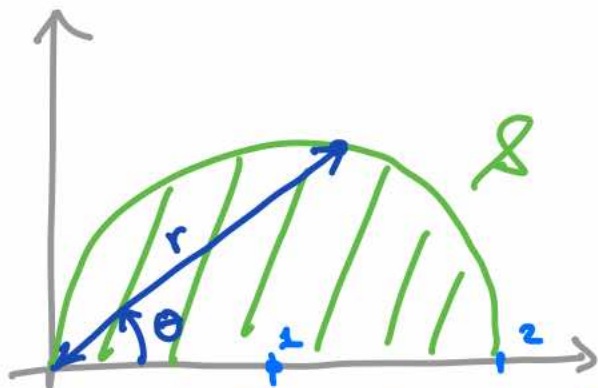
Convert it to polar coordinates:

$$\iint_{\mathcal{D}} r dA = \int_0^{\pi/2} \left( \int_0^{2\cos\theta} r^2 dr \right) d\theta$$

$$= \frac{8}{3} \int_0^{\pi/2} \cos^3 \theta d\theta$$

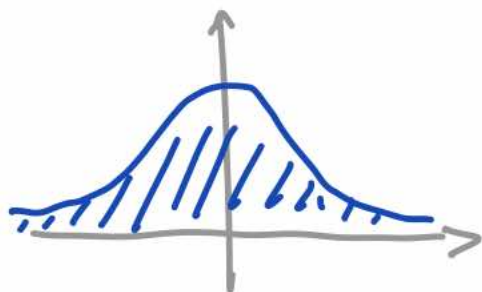
$$= \frac{8}{3} \int_0^{\pi/2} \cos \theta (1 - \sin^2 \theta) d\theta$$

$$= \frac{8}{3} \left[ \sin \theta - \frac{\sin^3 \theta}{3} \right]_0^{\pi/2} = \frac{8}{3} \left( 1 - \frac{1}{3} \right) = \frac{16}{9} \quad //$$



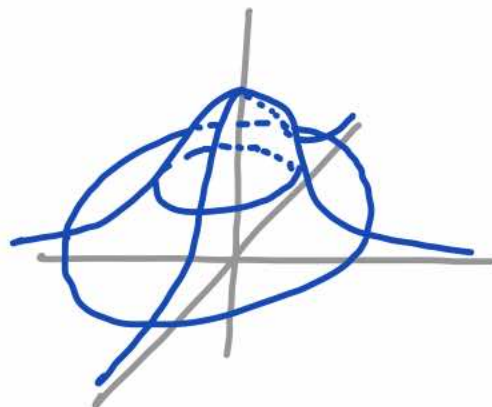
$$\begin{aligned} y &= \sqrt{2x-x^2} \rightarrow y^2 = 2x-x^2 \\ &\rightarrow (x^2-1)^2 + y^2 = 1 \text{ or } x^2 + y^2 = 2x \\ &\rightarrow r^2 = 2r \cos \theta \rightarrow \underline{r = 2 \cos \theta} \end{aligned}$$

Ex 4 / Compute  $I = \int_{-\infty}^{\infty} e^{-x^2} dx$ .



Consider the volume  $V$  under  $e^{-x^2-y^2}$ :

We will compute this in two ways.



$$\begin{aligned}
 \textcircled{1} \quad V &= \lim_{b \rightarrow \infty} \int_{-b}^b \int_{-b}^b e^{-x^2-y^2} dy dx \\
 &= \lim_{b \rightarrow \infty} \left( \int_{-b}^b e^{-y^2} dy \right) \left( \int_{-b}^b e^{-x^2} dx \right) \\
 &= \left( \lim_{b \rightarrow \infty} \int_{-b}^b e^{-x^2} dx \right)^2 = I^2
 \end{aligned}$$

$$\begin{aligned}
 \textcircled{2} \quad V &= \lim_{a \rightarrow \infty} \int_0^{2\pi} \int_0^a e^{-r^2} r dr d\theta \\
 &= \lim_{a \rightarrow \infty} \int_0^{2\pi} \left[ -\frac{1}{2} e^{-r^2} \right]_{r=0}^a d\theta \\
 &= \lim_{a \rightarrow \infty} \frac{1}{2} \int_0^{2\pi} (1 - e^{-a^2}) d\theta \\
 &= \lim_{a \rightarrow \infty} \pi (1 - e^{-a^2}) = \pi .
 \end{aligned}$$

$$\text{So } I^2 = \pi \quad \Rightarrow \quad I = \sqrt{\pi} .$$

